

Solution to Physics 108 Final (2014)

1- (1) From

$$\frac{n_{oil}}{s_o} + \frac{n_{oil}}{s_i} = (n_g - n_{oil}) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

or

$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{n_g - n_{oil}}{n_{oil}} \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

We have the inverse of the focal length f

$$\frac{1}{f} = \frac{n_g - n_{oil}}{n_{oil}} \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

Since $R_1 = \infty$, $R_2 = -10 \text{ cm}$, $n_g = 1.5$, $n_{oil} = 2.0$

$$\frac{1}{f} = \frac{1.5 - 2.0}{2.0} \cdot (-) \frac{1}{(-10 \text{ cm})} = -\frac{1}{40 \text{ cm}}$$

$$\therefore f = -40 \text{ cm}^*$$

1- (2) From $\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f}$, $s_i = \frac{s_o f}{s_o - f} = -\frac{40}{3} \text{ cm}$

$$y_i = \left(-\frac{s_i}{s_o} \right) \cdot y_o = \frac{2}{3} \text{ cm}^*$$

2-(a) With refractive index of the air = 1,

$$\begin{aligned}
 \begin{pmatrix} A & B \\ C & D \end{pmatrix} &= \begin{pmatrix} 1 & 0 \\ -\frac{1}{f_2} & 1 \end{pmatrix} \begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{f_1} & 1 \end{pmatrix} \\
 &= \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{3} & 1 \end{pmatrix} \\
 &= \begin{pmatrix} 1 & 2 \\ 1 & 3 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{3} & 1 \end{pmatrix} \\
 &= \begin{pmatrix} \frac{1}{3} & 2 \\ 0 & 3 \end{pmatrix}
 \end{aligned}$$

2-(2) $p = -\frac{D}{C} = -\infty \Rightarrow F_1$ at infinity

$q = -\frac{A}{C} = -\infty \Rightarrow F_2$ at infinity.

2-(3) $\begin{pmatrix} p_f \\ q_f \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} p_o \\ q_o \end{pmatrix} = \begin{pmatrix} \frac{1}{3} & 2 \\ 0 & 3 \end{pmatrix} \begin{pmatrix} p_o \\ q_o \end{pmatrix} = \begin{pmatrix} \frac{p_o}{3} \\ 0 \end{pmatrix}$

3-(1) Near normal incidence, the neighboring fringes correspond to a height (thickness) change δd given by

$$\frac{4\pi}{\lambda_0} n_{air} \delta d = 2\pi$$

$$\therefore \delta d \Big|_{\text{neighboring fringes}} = \frac{\lambda_0}{2}$$

With $\lambda_1 = 6000 \text{ \AA}$, the observation of 30 fringes means that the total thickness change is 30 times δd

$$\begin{aligned} \Delta d &= 30 \times \delta d \Big|_{\text{neighboring fringes @ } \lambda_1} \\ &= 30 \times \frac{\lambda_1}{2} = 90,000 \text{ \AA} \end{aligned}$$

3-(2) Since $\delta d \Big|_{\text{neighboring fringes @ } \lambda_2 = 4500 \text{ \AA}} = \frac{\lambda_2}{2} = 2250$

the number of "blue" fringes will be

$$N = \frac{\Delta d}{\delta d \Big|_{\text{blue}}} = \frac{90,000 \text{ \AA}}{2250 \text{ \AA}} = 40 \quad *$$

4-(1) Due to diffraction effect, the diameter of the CO_2 laser beam on the Moon should be roughly

$$\begin{aligned}d(\text{at Moon}) &\approx \left(\frac{2\lambda}{d}\right) \cdot L \\&= \frac{2 \times 10^{-5} \text{ meters}}{10^{-1} \text{ meters}} \times 3.84 \times 10^8 \text{ meters} \\&= 7.68 \times 10^4 \text{ meters} \# \end{aligned}$$

4-(2) The intensity of the laser beam at the Moon surface can be determined from energy conservation

$$\begin{aligned}(\pi/4) d^2 \cdot I_0 &\approx (\pi/4) d^2(\text{at Moon}) \cdot I(\text{at Moon}) \\ \therefore I(\text{at Moon}) &\approx I_0 \cdot \frac{d^2}{d^2(\text{at Moon})} \\ &= I_0 \cdot \left(\frac{d}{2\lambda L}\right)^2 = (1.7 \times 10^{-12}) I_0 \end{aligned}$$

5-(1) On the glass side

$$\tan \theta_{1B} = \frac{n_{air}}{n_g} = \frac{1}{2}$$

$$\therefore \theta_{1B} = 26.57^\circ$$

5-(2) From

$$V_s(\theta_{1B}) = \frac{n_g \cos \theta_{1B} - n_{air} \cos \theta_{2B}}{n_g \cos \theta_{1B} + n_{air} \cos \theta_{2B}}$$

$$= \frac{n_g - n_{air} \cdot \tan \theta_{1B}}{n_g + n_{air} \cdot \tan \theta_{1B}} \quad (\cos \theta_{2B} = \sin \theta_{1B})$$

$$= \frac{n_g^2 - n_{air}^2}{n_g^2 + n_{air}^2}$$

$$= \frac{3}{5}$$

$$R_s(\theta_{1B}) = |V_s(\theta_{1B})|^2 = \frac{9}{25} = 0.36 \#$$

5-3) At $\theta_1 = 45^\circ$,

$$\cos \tilde{\theta}_2 = \sqrt{1 - n_g^2 \cdot \sin^2 \theta_1} / n_{air}$$

$$= i$$

$$t_p = \frac{E_p^{(air)}}{E_{p,inc}^{(g)}} = \frac{2 n_g \cos \theta_1}{n_g \cos \tilde{\theta}_2 + n_{air} \cos \theta_1}$$

$$= \frac{2\sqrt{2}}{2i + \sqrt{2}/2}$$

$$\therefore |E_p^{(air)}| = |E_{p,inc}| \cdot \left| \frac{2\sqrt{2}}{\sqrt{2}/2 + i2} \right|$$

$$= \frac{4}{3} |E_{p,inc}|$$

✱

$$6-(a) \begin{pmatrix} -i \\ +i \end{pmatrix} = (-i) \begin{pmatrix} 1 \\ -1 \end{pmatrix}, \text{ linearly polarized} \\ \text{with } \alpha = -45^\circ$$

$$6-(b) \begin{pmatrix} -1+i \\ 1+i \end{pmatrix} = (-1+i) \begin{pmatrix} 1 \\ \frac{1+i}{-1+i} \end{pmatrix} = (-1+i) \begin{pmatrix} 1 \\ -i \end{pmatrix},$$

Right-circularly polarized light.

$$6-(c) \begin{pmatrix} -1-i \\ 1+i \end{pmatrix} = (-1-i) \begin{pmatrix} 1 \\ -1 \end{pmatrix}, \text{ linearly polarized} \\ \text{with } \alpha = -45^\circ$$

$$6-(d) \begin{pmatrix} i \\ 4 \end{pmatrix} = i \begin{pmatrix} 1 \\ -4i \end{pmatrix}, \text{ elliptically polarized} \quad \#$$

7-(c)

$$\vec{E}_{inc} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} = \begin{pmatrix} \cos 45^\circ \\ \sin 45^\circ \end{pmatrix}$$

$$M_{QWP}(\text{OAI} \parallel x\text{-axis}) = \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix}$$

$$\vec{E}_{out} = M_{QWP}(\text{OAI} \parallel x\text{-axis}) \vec{E}_{inc}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix} \begin{pmatrix} \frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

\therefore It is a right-circularly polarized light.

7-(2) Let the Jones matrix for a quarter-wave plate with OA || x-axis be

$$M_{\text{QWP}}(\text{OA} \parallel \hat{x}) = \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix}$$

then the Jones matrix for a quarter-wave plate with OA || y-axis will be

$$M_{\text{QWP}}(\text{OA} \parallel \hat{y}) = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$$

The total system Jones matrix is then given by

$$\begin{aligned} M &= M_{\text{QWP}}(\text{OA} \parallel \hat{y}) M_{\text{linear}}(\theta = 45^\circ) M_{\text{QWP}}(\text{OA} \parallel \hat{x}) \\ &= \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix} \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & i/2 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & -i/2 \\ i/2 & \frac{1}{2} \end{pmatrix} \end{aligned}$$

For an incident left-circular polarized light

$$\vec{E}_{inc}^{\sim} = \vec{E}_{left}^{\sim} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$$

After the polarization device combination,

$$\vec{E}_{out}^{\sim} = M \vec{E}_{inc}^{\sim}$$

$$= \begin{pmatrix} 1/2 & -i/2 \\ i/2 & 1/2 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix}$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} = \vec{E}_{inc}^{\sim}$$

\therefore It passes \vec{E}_{left}^{\sim} unscathed.

$$\text{If } \vec{E}_{inc}^{\sim} = \vec{E}_{right}^{\sim} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

$$\vec{E}_{out}^{\sim} = M \vec{E}_{right}^{\sim} = \begin{pmatrix} 1/2 & -i/2 \\ i/2 & 1/2 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} \\ -i/\sqrt{2} \end{pmatrix} = 0$$

\therefore It ~~also~~ removes \vec{E}_{right}^{\sim} completely.